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H.7000 STEP

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TITLE:

The electroacoustic effect with h-f ultrasonics

SOURCE:

Kishinev. Universitet. Uchenyye zapiski. v. 49, 1961, 26-28

TEXT: On the basis of quantum perturbation theory the electron-phonon interaction is studied for a system having an electron collision frequency smaller than the sonic frequency. Using a method of Herring and Vogt (Phys. Rev. 101, 944, 1956), the zero-current field strength inside an isolated semiconductor is derived as

$$E_{u}^{(j)} = \frac{2\pi g C_{e',p(l)} W_{o}kT}{\frac{1}{e^{m}}} \left[\frac{2\hbar^{\omega}}{\rho^{2}} \left(ch \frac{\hbar^{\omega}}{KT} - 1 \right) + sk \frac{\hbar^{\omega}}{KT} \right] \times \\ \times m_{e'} \int f_{o} \tau_{e'} d\epsilon / \int f_{o} \tau_{e'} \xi^{4} d\xi$$

$$(1.2)$$

. Card 1/3

from the equation of motion of electrons in the sound field.

$$C_{\mathbf{a}'p(l)} = \frac{q_{\mathbf{a}'}}{\left(\sum q_s^2 \frac{m_{\mathbf{a}'}}{m\beta'}\right)^{\frac{1}{2}}},$$

(1.2) holds for a nondegenerate carrier gas, whereas for a degenerate gas

$$E_a^{(l)} = -\frac{3 C_{a'p(l)} W_o(l) m_{a'}^2}{e} \cdot \frac{\hbar \omega}{\sqrt{\mu_o}}$$
 (1.4)

is valid. For high temperatures, ha/KT «1,

$$E_{a'} = -\frac{3Ca'p(l)}{2e} W_o^{(l)} kTm_{a'}^2 \left[\frac{\hbar\omega}{p^2} \left(\frac{\hbar\omega}{kT} \right)^2 + \frac{\hbar\omega}{KT} \right] \times \\ \times \int \exp(-e/KT) \tau_a de \int \exp(-e/KT) \tau_a t^4 dt.$$
(1.3)

Card 2/3

 $E_{\rm metal}/E_{\rm semic.} \simeq (kT/\mu_0)^{3/2}$. These relations are used to calculate eE for n-type Si and Ge for longitudinal and transverse oscillations. The electroacoustic effect can be used to determine both interaction constants, whereas other effects, as e. g. the piezoeffect, yield only one.